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# Stable isotopes in leaf water of terrestrial plants

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#### 55 ABSTRACT

56 Leaf water contains naturally occurring stable isotopes of oxygen and hydrogen in 57 abundances that vary spatially and temporally. When sufficiently understood, these can be 58 harnessed for a wide range of applications. Here, we review the current state of knowledge 59 of stable isotope enrichment of leaf water, and its relevance for isotopic signals incorporated into plant organic matter and atmospheric gases. Models describing evaporative enrichment 60 61 of leaf water have become increasingly complex over time, reflecting enhanced spatial and 62 temporal resolution. We recommend that practitioners choose a model with a level of 63 complexity suited to their application, and provide guidance. At the same time, there exists 64 some lingering uncertainty about the biophysical processes relevant to patterns of isotopic 65 enrichment in leaf water. An important goal for future research is to link observed variations 66 in isotopic composition to specific anatomical and physiological features of leaves that reflect 67 differences in hydraulic design. New measurement techniques are developing rapidly, enabling determinations of both transpired and leaf water  $\delta^{18}$ O and  $\delta^{2}$ H to be made more 68 easily and at higher temporal resolution than previously possible. We expect these 69 70 technological advances to spur new developments in our understanding of patterns of stable 71 isotope fractionation in leaf water. 72

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2

#### 74 INTRODUCTION

75 In this review, we focus on how stable isotope ratios of oxygen and hydrogen vary in 76 leaf water. The stable isotope composition of leaf water significantly influences isotopic 77 signatures of a number of important biological and atmospheric processes. For example, the 78 oxygen isotope composition of leaf water partly controls the oxygen isotope compositions of atmospheric CO<sub>2</sub> (Farquhar et al. 1993; Cuntz et al. 2003; Welp et al. 2011) and atmospheric 79 80 O<sub>2</sub> (Dole et al. 1954; Hoffmann et al. 2004; Luz & Barkan 2011). Sugars and other 81 metabolites formed in leaves incorporate the leaf water isotopic signal, which is then retained 82 in structural organic compounds, such as cellulose (Saurer, Aellen & Siegwolf 1997; Roden, 83 Lin & Ehleringer 2000; Barbour 2007; Gessler et al. 2014). The leaf water signal is also 84 preserved in leaf waxes (Smith & Freeman 2006; Sachse et al. 2010; Kahmen et al. 2013a; 85 Kahmen, Schefuss & Sachse 2013b), components of which can persist in the environment for 86 millions of years (Eglinton & Eglinton 2008). Thus, leaf water derived isotopic signals can 87 be useful for constraining models of the global carbon cycle, reconstructing past climates, 88 retrospectively analysing plant physiological responses to the environment, and for assigning 89 geographic origins to plant materials and plant-derived products (Dawson et al. 2002; West et 90 al. 2006b). All of these various applications rely on a firm understanding of the mechanisms 91 that control leaf water isotopic enrichment.

92 For plant water, isotopic abundances are generally expressed relative to the
93 international standard VSMOW (Vienna Standard Mean Ocean Water) (Coplen 2011). This
94 is accomplished using δ notation:

95 
$$\delta_p = \frac{R_p - R_{Std}}{R_{Std}},$$
 (1)

96 where  $R_p$  is the isotope ratio (*e.g.*, <sup>18</sup>O/<sup>16</sup>O or <sup>2</sup>H/<sup>1</sup>H) of a plant water sample and  $R_{Std}$  is that 97 of the standard. The resulting  $\delta$  values are typically multiplied by 1000, so that the relative 98 deviation of the isotope ratio of the sample from that of the standard is expressed as per mil 99 (‰).

100 The stable isotope composition of plant water is influenced firstly by the plant's 101 source water; this is mainly water taken up by roots from the soil. Soil water for terrestrial 102 plants generally derives from local precipitation. The stable isotope composition of 103 precipitation can vary both geographically and temporally. The  $\delta^{18}$ O and  $\delta^{2}$ H of precipitation 104 have been shown to vary in conjunction with temperature, altitude, latitude, distance from the 105 coast, and with the amount of precipitation falling in a given event (Rozanski, Araguas-106 Araguas & Gonfiantini 1993; Araguas-Araguas, Froehlich & Rozanski 2000; Bowen 2010; 107 Munksgaard *et al.* 2012). A representation of geographic variation in the  $\delta^{18}$ O and  $\delta^{2}$ H of 108 mean annual precipitation across the global land surface is shown in Figures 1A and 1B. 109 Any given precipitation event will mix into an existing soil water pool. Thus,

although there can exist relatively large variation in  $\delta^{18}$ O and  $\delta^{2}$ H from one precipitation 110 event to the next (Munksgaard et al. 2012; Munksgaard et al. 2015), the soil water pool being 111 112 accessed by plants will likely be buffered to some extent against these short term variations. 113 The extent to which the isotopic composition of the soil water pool can be linked to 114 individual precipitation events is an area of current interest (Tang & Feng 2001; Brooks et al. 115 2010; Thomas et al. 2013; Gessler et al. 2014). It is particularly relevant, for example, to the detection of long-term records of tropical cyclone activity in tree rings (Miller et al. 2006), 116 117 because tropical cyclones are predominantly associated with isotopically light precipitation 118 (Gedzelman & Arnold 1994; Lawrence & Gedzelman 1996; Munksgaard et al. 2015). The 119 isotopic composition of soil water can also be affected by evaporation. This causes the soil water near the evaporating front to become enriched in <sup>18</sup>O and <sup>2</sup>H compared to the soil water 120 121 at depth (Allison, Barnes & Hughes 1983; Barnes & Allison 1983).

122 For the most part, the isotopic composition of water in non-transpiring plant organs 123 (*i.e.*, roots, stems, etc.) has been shown to match that of the water available to the plant in the 124 soil, indicating that there is little to no stable isotope fractionation associated with absorption 125 of water by roots and transport in xylem (White 1989; Ehleringer & Dawson 1992; Dawson 1993). The exception to this rule is that  $\delta^2 H$  has been observed to shift with water uptake 126 127 and/or transport in salt tolerant coastal plants (Lin & Sternberg 1993) and phreatophytic desert shrubs (Ellsworth & Williams 2007). No simultaneous  $\delta^{18}$ O fractionation was 128 observed, indicating that the cause of the isotope effect was specific only to hydrogen 129 130 isotopes in water and not oxygen isotopes, or that the isotope effect for oxygen was too small 131 to be detected.

Transpiration results in isotopic enrichment at the sites of evaporation within leaves. The isotopically enriched water can then diffuse away from the evaporative sites into other parts of the leaf. The resulting bulk leaf water enrichment generally shows a diurnal pattern, with a daily maximum in the early afternoon associated with the minimum daily relative humidity, and a daily minimum in the early morning reflecting a progressive relaxation of the enrichment through the night (Figure 2). Enriched leaf water can also be transported in the phloem to developing sink organs such as seeds (Figure 2). In the following sections, we examine in detail the environmental and physiological controls over the stable isotope

140 enrichment of the evaporative sites and the bulk leaf water.

141

#### 142 **LEAF WATER**

#### 143 **Evaporative sites**

Isotopic enrichment of leaf water as a result of the evaporative process of 144 145 transpiration was first observed by Gonfiantini et al. (1965). In the same year, a model for 146 predicting the isotopic enrichment that should take place at the surface of an evaporating 147 body of water was published by Craig and Gordon (1965). This model can be applied to the 148 isotopic composition of water at the evaporative sites within leaves ( $\delta_e$ ). Here, it is 149 convenient to express the isotopic composition of the evaporative sites as enrichment in the 150 heavier isotopes compared to source water, to account for the influence of different source-151 water isotopic signatures among plants. The enrichment of any plant water sample above 152 source water  $(\Delta_p)$  can be expressed as

153 
$$\Delta_p = \frac{\delta_p - \delta_s}{1 + \delta_s} , \qquad (2)$$

where  $\delta_p$  is the  $\delta$  value of the plant water sample and  $\delta_s$  is that of source water. Here again,  $\Delta_p$ ,  $\delta_p$ , and  $\delta_s$  are often expressed as per mil. If this is the case, the  $\delta_s$  in the denominator on the right side of the equation must be divided by 1000. A list of the main symbols and abbreviations used throughout the text is given in Table 1.

The Craig-Gordon model, as modified for application to leaves by subsequent authors
(Dongmann *et al.* 1974; Flanagan, Comstock & Ehleringer 1991; Farquhar & Lloyd 1993),
can be approximated by

161

$$\Delta_e \approx \varepsilon^+ + \varepsilon_k + (\Delta_v - \varepsilon_k) \frac{w_a}{w_i},\tag{3}$$

where  $\Delta_e$  is the enrichment of evaporative site water above source water,  $\epsilon^+$  is the equilibrium 162 163 fractionation between liquid water and vapour,  $\varepsilon_k$  is the kinetic fractionation for combined 164 diffusion through the stomata and the boundary layer,  $\Delta_v$  is the isotopic enrichment of 165 atmospheric vapour compared to source water, and  $w_a/w_i$  is the ratio of the water vapour mole 166 fraction in the air relative to that in the intercellular air spaces. Thus,  $w_a/w_i$  is the relative 167 humidity, but with the saturation water vapour mole fraction in the denominator calculated 168 for leaf temperature rather than air temperature. If leaf temperature and air temperature are 169 equal,  $w_a/w_i$  is exactly equal to the relative humidity of the air. The  $\Delta_v$  is calculated with 170 respect to source water as shown in Eqn 2, and it typically has a negative value due to the

equilibrium isotope effect between liquid and vapour. Equation 3 is a convenient

approximation for the precise form of the model, as given by (Farquhar, Cernusak & Barnes2007)

$$\Delta_e = (1 + \varepsilon^+) \left[ (1 + \varepsilon_k) \left( 1 - \frac{w_a}{w_i} \right) + \frac{w_a}{w_i} (1 + \Delta_v) \right] - 1.$$
(4)

175 Note that in Eqn 3, the calculation can be readily performed with all isotopic terms expressed 176 as per mil, whereas for Eqn 4, it is more straightforward to make the calculation with the 177 isotopic terms not expressed as per mil, and then to multiply the result by 1000 afterward to 178 return to per mil notation. In the supplementary material, we provide a Microsoft Excel 179 spreadsheet with a combined data set of leaf water observations, which also contains a 180 worked example of how to perform the calculation shown in Eqn 4. The difference between 181  $\Delta_{\rm e}$  calculated with Eqn 3 and that calculated with Eqn 4 is small for oxygen, on the order of 182 0.1‰. For hydrogen, it is larger, on the order of 1 to 2‰.

183 The equations shown above for predicting  $\Delta_e$  assume isotopic steady state. Isotopic 184 steady state means that the isotopic composition of the the transpired water vapour is equal to 185 that of the source water supplying the leaf (Craig & Gordon 1965; Harwood et al. 1998; 186 Farquhar & Cernusak 2005). The condition of non-steady state, when the transpired water 187 has an isotopic composition differing from that of source water, will be discussed below. In 188 general, it has been observed that leaf water enrichment tends to be near to steady state 189 during the day in leaves that have relatively open stomata and do not show a high degree of 190 succulence (Cernusak et al. 2008).

The equilibrium fractionation varies as a function of temperature (Bottinga & Craig
1969; Majoube 1971; Horita & Wesolowski 1994). It can be calculated according to the
following equations (Majoube 1971), with that for <sup>18</sup>O shown first, followed by that for <sup>2</sup>H:

194 
$$\varepsilon_0^+(\%_0) = \left[ e^{\left(\frac{1.137}{(273+T)^2} \times 10^3 - \frac{0.4156}{273+T} - 2.0667 \times 10^{-3}\right)} - 1 \right] \times 1000 ,$$

195 
$$\varepsilon_{H}^{+}(\%_{0}) = \left[ e^{\left(\frac{24.844}{(273+T)^{2}} \times 10^{3} - \frac{76.248}{273+T} + 52.612 \times 10^{-3}\right)} - 1 \right] \times 1000 .$$
 (6)

196 The right sides of the equations have been multiplied by 1000, so that  $\varepsilon^+$  is here expressed as 197 per mil. The symbol  $\varepsilon_0^+(\infty)$  in Eqn 5 refers to the isotope fractionation for <sup>18</sup>O, and  $\varepsilon_H^+(\infty)$  in 198 Eqn 6 refers to that for <sup>2</sup>H. The *T* in these equations refers to the leaf temperature in degrees 199 Celsius. The  $\varepsilon_k$  in Eqns 3 and 4 can be calculated as (Farquhar *et al.* 1989)

200 
$$\varepsilon_k^O(\%_0) = \frac{28r_s + 19r_b}{r_s + r_b},$$
 (7)

201 
$$\varepsilon_k^H(\%_0) = \frac{25r_s + 17r_b}{r_s + r_b}.$$
 (8)

(5)

The  $\varepsilon_k^O(\%_0)$  is  $\varepsilon_k$  for <sup>18</sup>O expressed as per mil, and  $\varepsilon_k^H(\%_0)$  is the same for <sup>2</sup>H. The  $r_s$  and  $r_b$ 202 in Eqns 7 and 8 are the stomatal and boundary layer resistances, respectively (m<sup>2</sup> s mol<sup>-1</sup>); 203 204 they are the inverses of the stomatal and boundary layer conductances. The 28 and 19 in Eqn 7 are fractionation factors for diffusion of water molecules containing <sup>18</sup>O through the 205 206 stomata and boundary layer, expressed as per mil. The values 25 and 17 in Eqn 8 are those 207 same fractionation factors for <sup>2</sup>H (Merlivat 1978). It has been suggested that these values 208 should be revised (Cappa et al. 2003). However, subsequent measurements indicated that the 209 fractionation factors originally assigned are the more correct values (Luz et al. 2009).

210 If the water vapour in the air is in isotopic equilibrium with source water, then  $\Delta_v$  will 211 approximately equal  $-\epsilon^+$ . In that case, Eqn 3 will condense to

$$\Delta_e \approx \left(\varepsilon^+ + \varepsilon_k\right) \left(1 - \frac{w_a}{w_i}\right). \tag{9}$$

Equation 9 demonstrates the strong role that the relative humidity term  $w_a/w_i$  plays in determining the isotopic enrichment of leaf water at the sites of evaporation.

212

Figure 3 shows the relationships between observed daytime bulk leaf water isotopic enrichment and the air relative humidity and Craig-Gordon predictions for a large dataset collected under natural field conditions across a sub-continental rainfall gradient in northern Australia (Kahmen *et al.* 2013a). The analysis shows both the importance of the relative humidity term in driving daytime leaf water stable isotope enrichment (Figures 3A and 3C), and that the Craig-Gordon equation captures much of the observed variation across a largescale environmental gradient (Figures 3B and 3D).

This analysis also highlights an important difference between <sup>18</sup>O and <sup>2</sup>H. For <sup>18</sup>O, 222 223 the air relative humidity predicts nearly as much variation in the observed leaf water enrichment as does the full Craig-Gordon model, with  $R^2$  of 0.78 for the former versus 0.86 224 for the latter (Figures 3A and 3B). For <sup>2</sup>H, on the other hand, the air relative humidity 225 226 predicts only a little more than half the variation predicted by the full Craig-Gordon model, with  $R^2$  of 0.52 for the former versus 0.92 for the latter (Figures 3C and 3D). This 227 228 demonstrates the importance of the isotopic disequilibrium between air vapour and source water for predicting  $\Delta_e$  for <sup>2</sup>H in comparison to <sup>18</sup>O. This disequilibrium can be expressed as 229  $\epsilon^++\Delta_v$ . The contrast between <sup>2</sup>H and <sup>18</sup>O in the sensitivity of  $\Delta_e$  to  $\epsilon^++\Delta_v$  comes about 230 because, for <sup>18</sup>O,  $\epsilon^+$ ,  $\Delta_v$ , and the difference between them are typically small in absolute 231 value compared to  $\varepsilon_k$ , whereas the opposite is true for <sup>2</sup>H. For <sup>2</sup>H, the disequilibrium term 232  $\varepsilon^+ + \Delta_v$  can easily be larger than  $\varepsilon_k$  in absolute value, with either positive or negative values 233 possible. Thus, the predicted  $\Delta_e$  for <sup>18</sup>O is dominated by the kinetic fractionation,  $\varepsilon_k$ ; whereas 234

for <sup>2</sup>H, the predicted  $\Delta_e$  is dominated by the equilibrium fractionation,  $\epsilon^+$ , and by the air vapour disequilibrium term,  $\epsilon^+ + \Delta_v$ .

237 The role of the atmospheric vapour isotopic composition in controlling  $\Delta_e$  can be 238 further appreciated by examining the limiting case where relative humidity is saturated, such that  $w_a/w_i=1$ . In this case, Eqn 3 reduces to  $\varepsilon^++\Delta_v$ ; and, the isotopic disequilibrium between 239 air vapour and source water then controls  $\Delta_{e}$ . While this limiting scenario usually only 240 241 occurs at night, it emphasises the importance of atmospheric vapour in influencing leaf water 242 enrichment (Farquhar & Cernusak 2005; Helliker & Griffiths 2007), as well as the general 243 importance of having a reasonably accurate estimate of  $\Delta_v$  for predicting  $\Delta_e$ , especially with respect to <sup>2</sup>H. In humid-zone epiphytes that use Crassulacean acid metabolism, this 244 245 phenomenon creates an opportunity to reconstruct the isotope ratio of atmospheric water 246 vapour from the epiphyte's organic matter (Helliker 2014).

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#### 248 Bulk leaf water

The term 'bulk leaf water' generally refers to a water sample obtained by extraction from a whole leaf. A bulk leaf water sample may or may not contain the water of the major veins, depending on the sampling protocol of the individual researcher. Leaf water excluding the major veins has also been referred to as 'lamina leaf water'. It is important to note that in the vast majority of plants, such a sample will also contain water associated with minor veins. Here we use  $\delta_L$  to refer to the isotopic composition of bulk leaf water, and  $\Delta_L$  to refer to its enrichment above source water.

256 Early measurements indicated that the Craig-Gordon model tended to overestimate  $\Delta_{\rm L}$ 257 (Allison, Gat & Leaney 1985; Leaney et al. 1985; Bariac et al. 1989; Walker et al. 1989; 258 Yakir, DeNiro & Gat 1990; Flanagan et al. 1991; Walker & Lance 1991). To illustrate this 259 phenomenon, we compiled leaf water isotopic data from a number of published datasets, 260 along with the Craig-Gordon prediction of leaf water enrichment corresponding to each 261 observation (Supplementary material). The dataset contains 118 species, sampled across a 262 range of tropical and temperate sites from both northern and southern hemispheres. It is 263 limited to daytime observations of C<sub>3</sub> plants under natural field conditions. Figure 4 presents 264 the results for the proportional difference between the predicted Craig-Gordon enrichment and the observed bulk leaf water enrichment  $(1-\Delta_I/\Delta_e)$ . The analysis confirms that observed 265  $1-\Delta_{\rm I}/\Delta_{\rm e}$  is larger than zero for both <sup>18</sup>O (P<0.001; n=722) and <sup>2</sup>H (P<0.001; n=362), with 266 average proportional differences 0.12 for  $^{18}$ O and 0.24 for  $^{2}$ H. 267

268 The explanation for the generally lower observed value of  $\Delta_L$  compared to  $\Delta_e$  has attracted considerable research effort, because it is important to determine which leaf water 269 270 signal is most relevant to the various applications that depend upon it. Two models have 271 been proposed to explain this pattern when steady state conditions can reasonably be 272 expected: a two-pools model, based on two discrete pools of water within the leaf, with one 273 of them being unenriched xylem water (Leaney et al. 1985; Yakir, Deniro & Rundel 1989; 274 Yakir et al. 1990; Roden & Ehleringer 1999; Song et al. 2015a); and an advection-diffusion, 275 or Péclet, model (Farquhar & Lloyd 1993; Farquhar & Gan 2003; Barnes, Farquhar & Gan 276 2004).

277 If the two-pools model is assumed to comprise unenriched source water and enriched
278 evaporative site water, it can be written as (Leaney *et al.* 1985; Song *et al.* 2015a)

$$\Delta_L = (1-\phi)\Delta_e$$
,

where  $\phi$  is the proportion of leaf water that is unenriched xylem water, presumably residing mainly in the major veins and ground tissue associated with them. In this model, the overestimation of  $\Delta_L$  by the Craig-Gordon model is due to the contribution from the unenriched pool.

284 Rather than two discrete pools, the Péclet model describes gradients of enrichment 285 within the leaf water. In the Péclet model, advection of less enriched water by the 286 transpiration stream opposes the back-diffusion of isotopically enriched water from the 287 evaporative sites (Farquhar & Lloyd 1993). When advection overwhelms diffusion, the bulk 288 leaf water enrichment will be less than that predicted by the Craig-Gordon equation. 289 Accordingly, the proportional difference,  $1-\Delta_L/\Delta_e$ , is predicted to increase with increasing 290 transpiration rate. This particular feature is an important distinction between the Péclet 291 model and the two-pools model: the Péclet model predicts that the deviation of the bulk leaf 292 water from the Craig-Gordon predicted enrichment should record information about the 293 transpiration rate. In contrast, the two-pools model does not predict such an effect.

The Péclet number, which is dimensionless, represents the extent to which diffusion is overwhelmed by advective counter-flow (Ikeda 1983). It was originally developed to describe the ratio between convective and conductive heat transfer by Jean Claude Eugene Péclet, and has since been applied more generally to describe advection-diffusion effects on mass transport processes in permeable media. For leaves, the Péclet number,  $\wp$ , can be defined as *vl/D*, where *v* is the velocity of water movement (m s<sup>-1</sup>), *l* is the distance (m) from the evaporative sites over which the Péclet effect is occurring, and *D* is the diffusivity of the

(10)

heavy isotopologue in water (m<sup>2</sup> s<sup>-1</sup>). The *D* is temperature dependent, and can be modelled as a function of leaf temperature as (Cuntz *et al.* 2007),

$$D_0 = 97.5 \times 10^{-9} e^{\left(-\frac{577}{T-145}\right)},\tag{11}$$

$$D_H = 98.7 \times 10^{-9} e^{\left(-\frac{577}{T-145}\right)},$$
(12)

where  $D_0$  is the diffusivity for H<sub>2</sub><sup>18</sup>O,  $D_H$  is that for H<sup>2</sup>HO, and T is leaf temperature in °C. 305 The velocity of advection can further be described as kE/C, where E is the transpiration rate 306 (mol m<sup>-2</sup> s<sup>-1</sup>), C is the molar concentration of water ( $5.55 \times 10^4$  mol m<sup>-3</sup>), and k is a scaling 307 factor to account for the tortuosity of the water path. The term E/C gives the velocity as if 308 309 water were moving as a slab perpendicular to the leaf surface. The true velocity must be 310 faster than the slab velocity, because water moves in a tortuous path through the leaf. The 311 scaling factor k represents the ratio of the true velocity to the slab velocity. Combining the 312 above terms gives the following definition for the Péclet number:

313 
$$\wp = \frac{klE}{CD}.$$
 (13)

It is convenient to combine k and l into a single term, which has been called the effective path length, L (Farquhar & Lloyd 1993):

$$\beta = \frac{LE}{CD}.$$
(14)

Ignoring the water in veins for the moment, the Péclet model applied to the average leaf
lamina then predicts the following relationship with the evaporative site water enrichment
(Farquhar & Lloyd 1993):

320  $\Delta_L = \Delta_e \left( \frac{1 - e^{-\psi}}{\psi} \right). \tag{15}$ 

Equation 15 indicates that the smaller the Péclet number, the more similar  $\Delta_L$  will be to  $\Delta_e$ . It also predicts a continuous isotopic gradient from the sites of evaporation to the source water, modelled as an exponential decay along a cylindrical flow path.

324 Because Ø includes terms for both the effective path length and the transpiration rate, 325 a change in either one is predicted to alter the relationship between  $\Delta_L$  and  $\Delta_e$ . This is shown 326 schematically in Figure 5. In practice, values for the effective path length L have been 327 difficult to determine directly. Thus, they have generally been fitted using Eqns 14 and 15, 328 which therefore involves comparing observed bulk leaf water enrichment with the predicted 329 Craig-Gordon enrichment. One consideration that could lead to biased estimates of L when 330 fitted in this way is that unenriched vein water could also contribute to the difference between  $\Delta_{\rm L}$  and  $\Delta_{\rm e}$ , as described by the two pool model. If this were the case, L would be 331 332 overestimated if unenriched vein water were not accounted for prior to fitting Eqn 15.

333 Farquhar and Gan (2003) improved upon the one dimensional Péclet model described 334 above by separating Péclet effects in the leaf xylem and lamina. In practical terms, this 335 provides a means of combining the two-pool concept of dilution of leaf water enrichment by 336 relatively unenriched vein water with the lamina Péclet model. Furthermore, it allows the 337 advection-diffusion behaviour to be expressed in vein water as well as in mesophyll water. 338 Such a consideration is important because some observations suggest that vein water can become enriched in <sup>18</sup>O (Gan *et al.* 2002; Gan *et al.* 2003). Farquhar and Gan (2003) 339 340 suggested that most vein water should be found in the major veins and associated ground 341 tissue, with the proportion of leaf water in higher order minor veins being relatively small. If 342 the proportion of leaf water in minor veins were considered negligible, the bulk leaf water 343 enrichment could then be described as (Farquhar & Gan 2003),

$$\Delta_L = \Delta_e \left[ \phi_x e^{-\wp_r} + (1 - \phi_x) \frac{1 - e^{-\wp}}{\wp} \right],\tag{16}$$

where  $\phi_x$  is the proportion of leaf water in major veins, and  $\wp_r$  is the total radial Péclet number, equal to the sum of the lamina radial Péclet number,  $\wp$ , and the veinlet Péclet number,  $\wp_{rv}$ . Equation 16 allows for part of the difference between  $\Delta_L$  and  $\Delta_e$  to be accounted for by the relatively unenriched vein water. Fitting mesophyll effective path lengths with Eqn 16 should therefore provide more realistic estimates than with Eqn 15, but has the added complexity that values need to be assigned for  $\phi_x$  and  $\wp_r$  (e.g. Ripullone *et al.* 2008), or these need to be fitted simultaneously (e.g. Gan *et al.* 2003).

344

352 A convenient way to probe observed leaf water isotopic composition for evidence of 353 Péclet effects is to plot the proportional difference between  $\Delta_e$  and  $\Delta_L$  as a function of 354 transpiration rate. Such plots have yielded variable results, with some authors finding a 355 positive relationship, as predicted by the Péclet model (Barbour et al. 2000b; Ripullone et al. 356 2008; Loucos *et al.* 2015), and others, either no detectable relationship or a negative 357 relationship (Roden & Ehleringer 1999; Cernusak, Wong & Farquhar 2003; Song et al. 2013; 358 Roden et al. 2015; Song et al. 2015a). For our combined dataset given in the supplementary 359 material and shown in Figure 4, we find no relationship between  $1-\Delta_L/\Delta_e$  and transpiration rate for either  $\Delta^{18}$ O or  $\Delta^{2}$ H. 360

Given the conceptual realism in the Péclet model, it has been difficult to explain why in some cases there is no observable relationship between  $1-\Delta_{\rm L}/\Delta_{\rm e}$  and *E*. One explanation might be changes in the effective path length as transpiration rate varies (Kahmen *et al.* 2008; Song *et al.* 2013; Loucos *et al.* 2015). Water supply to the mesophyll is predominantly via the minor veins (Sack & Holbrook 2006). Once in the mesophyll, water movement to the 366 sites of evaporation can proceed through three parallel pathways: symplastic movement 367 through plasmodesmata, transcellular movement across cell membranes through aquaporins, 368 and apoplastic flow in cell walls that are not suberised (Steudle, Murrmann & Peterson 1993). 369 Although it can reasonably be expected that most flow will occur through the apoplast 370 (Brodribb, Feild & Jordan 2007), each of these pathways is nonetheless likely to be 371 associated with its own effective path length (Barbour & Farquhar 2004), and the possibility exists that the relative activity of these pathways may change with transpiration rate. In 372 373 addition, water may not always evaporate in the vicinity of the stomatal pore (Rockwell, 374 Holbrook & Stroock 2014; Buckley 2015), as is generally assumed. These considerations 375 have potential to obscure the positive relationship between  $1-\Delta_L/\Delta_e$  and E that is predicted by 376 the Péclet model, because changes to L could compensate for changes in E, thereby 377 decoupling  $\wp$  from *E* (Cernusak & Kahmen 2013; Song *et al.* 2013).

378 The idea that xylem can be variably coupled to the mesophyll to give distinct pools of 379 water of different volume and function within the leaf has been suggested in relation to leaf 380 hydraulics (Zwieniecki, Brodribb & Holbrook 2007; Canny et al. 2012), and would support 381 the idea of isotopic compartmentalisation of leaf water (Yakir et al. 1989; Yakir 1992; Yakir 382 et al. 1994). Looking at the rehydration kinetics of leaves of different species, Zwieniecki et 383 al. (2007) considered three observed patterns of hydraulic design: 1) where the vein is 384 hydraulically separated from the rest of the leaf; 2) where the epidermis is hydraulically 385 linked to the veins through the bundle sheath extension, but the mesophyll remains separated, 386 and 3) where all tissues are equally well coupled (Figure 6). Such compartmentalisation 387 could be created by both the internal organisation of leaf tissues, leading to variable degrees 388 of physical contact between different structures, and by the number, activity, and resistance 389 of the different pathways for water movement. A reasonable hypothesis, based on these 390 observations, is that different residence times will occur for different pools of water within 391 the leaf, introducing further variation into observed relationships between  $1-\Delta_L/\Delta_e$  and E, 392 because some pools of water would carry a memory of previous leaf water enrichment 393 conditions, whereas others would not.

Conifer needles fit within Design 1 of Figure 6, consisting of a singular vascular bundle surrounded by transfusion tissue and a thick-walled endodermis, which likely provides high radial resistance and physical separation between xylem and mesophyll. Consistent with this concept, it was recently observed that a two-pool model was sufficient to explain the difference between the Craig-Gordon prediction and the observed bulk leaf water 399 enrichment in two pine species (Roden et al. 2015). Water pools may also exist within the 400 mesophyll. In Eucalyptus pauciflora (snowgum) mesophyll cells shrank equally during 401 transpiration (Canny & Huang 2006), whereas in *Gossypium hirsutum* (cotton), cavity and 402 spongy mesophyll cells shrank more than matrix cells (Canny et al. 2012), suggesting that 403 different pools of water differentially supported evaporative demand. Leaf shrinkage of 404 tissues has also been linked more generally to the decline in extra-xylary hydraulic 405 conductance (Scoffoni et al. 2014), which could further contribute to hydraulic 406 compartmentalisation under conditions of water stress.

407

#### 408 **Progressive enrichment**

Sampling leaf tissue at a sub-leaf scale has revealed spatial patterns of isotopic
enrichment within leaves (Figure 7). Here, the isotopic composition tends to become
progressively enriched towards the tip of the leaf and out from the mid-vein (Bariac *et al.*1994; Wang & Yakir 1995; Helliker & Ehleringer 2000; Gan *et al.* 2002; Santrucek *et al.*2007).

414 This spatial pattern was initially explained using a string of lakes model, which 415 assumed a string of inter-connected pools of water within the leaf with differing isotope 416 compositions (Gat & Bowser 1991; Helliker & Ehleringer 2000; Helliker & Ehleringer 417 2002). Farquhar and Gan (2003) improved upon this model by including Péclet effects in 418 both mesophyll and veins (Figure 8). This enabled predictions of progressive enrichment of 419 xylem water in monocot leaves with distance from the base of the leaf. The predictions 420 matched relatively well the observed pattern in maize (Farquhar & Gan 2003; Gan et al. 421 2003). Ogée et al. (2007) then further improved upon this model by incorporating non-steady 422 state effects.

423 The progressive enrichment observed in both monocot and dicot leaves suggests that 424 back-diffusion occurs from the mesophyll back into the vein, allowing some evaporative 425 enrichment to be passed via the xylem from central and basal portions of the leaf to 426 downstream leaf sections. Such spatial variation can be described by three Péclet numbers 427 (Figure 8): 1) a radial Péclet number at the interface between xylem and mesophyll ( $\wp_{ry}$ ) 428 which allows for the leaf veinlet xylem water to become enriched above petiole water; 2) a 429 radial Péclet number associated with the mesophyll tissue ( $\wp$ ) which is likely to be small; 430 and 3) a longitudinal Péclet number ( $\beta_1$ ), allowing progressive enrichment of the xylem in 431 major veins in the direction of water movement, which is large, meaning mass transfer of

enrichment is mainly driven by advection. A derivation of this two dimensional Péclet model
with component Péclet numbers is given by Farquhar & Gan (2003).

In a general sense, observations of progressive enrichment provide strong support for the concept of Péclet effects in leaf water, because increasing enrichment of vein water with increasing distance from the midrib and the leaf base would not occur if some enriched water did not back diffuse from the evaporative sites into the veins, against the advective flow of the vein water.

439

#### 440 Transpired water

441 Water vapour leaving a leaf during transpiration originates directly from water at the 442 evaporative sites. Thus, it makes intuitive sense that the isotopic composition of transpired 443 water vapour ( $\delta_E$ ) should be related to that of evaporative site water ( $\delta_e$ ). When the Craig-444 Gordon model is written in a form that does not assume isotopic steady state, it predicts the 445 following relationship between  $\delta_e$  and  $\delta_E$ :

446

$$\delta_e \approx \delta_E + \varepsilon^+ + \varepsilon_k + (\delta_v - \delta_E - \varepsilon_k) \frac{w_a}{w_i}.$$
(17)

From Eqn 17, it can be seen that  $\delta_E$  is a necessary component for predicting  $\delta_e$  under nonsteady state conditions. When steady state is assumed,  $\delta_E$  is set equal to  $\delta_s$ . Making this substitution then leads to the widely used formulation shown in Eqn 3. If  $\delta_E$  is measured experimentally, Eqn 17 provides a useful means of estimating the isotopic composition of the evaporative sites under non-steady state conditions (Harwood *et al.* 1998).

452 The steady-state assumption of  $\delta_E$  being equal to  $\delta_s$  results from mass balance 453 constraints on leaf water dynamics, as shown in the following equation (Dongmann 1974; 454 Farquhar & Cernusak 2005):

$$\frac{d(W\delta_L)}{dt} = E(\delta_s - \delta_E).$$
(18)

Here *W* is the leaf water concentration and  $\delta_{\rm L}$  is the isotope composition of leaf water. The product of the two is termed isostorage (Farquhar & Cernusak 2005). The term  $E(\delta_{\rm s}-\delta_{\rm E})$ describes the difference between the isotopic flux of water molecules into  $(E\delta_{\rm s})$  and out of ( $E\delta_{\rm E}$ ) the leaf, and thus can be viewed as the net isoflux. Equation 18 states that the rate of change of leaf water isostorage is equal to the net isoflux of water into or out of the leaf. With the leaf at isotopic steady state, leaf water isostorage would be constant (*i.e.*,  $\frac{d(W\delta_{\rm L})}{dt} =$ 0). Accordingly the net isoflux would be zero, such that  $\delta_{\rm E}$  must be equal to  $\delta_{\rm s}$ .

463 Motivated by the need to address the conditions under which isotopic steady state occurs (*i.e.*,  $\delta_{\rm E} = \delta_{\rm s}$ ), several authors have used isotope ratio laser spectrometry coupled to a 464 gas exchange system to explore the variability of  $\delta_{\rm E}$  in response to environmental conditions 465 466 (Wang et al. 2012; Simonin et al. 2013; Dubbert et al. 2014; Song et al. 2015b). In a 467 laboratory study conducted on tobacco and citrus leaves, Simonin et al. (2013) observed that 468  $\delta_E$  was variable and deviated from  $\delta_s$  as long as instability was present in any of the 469 environmental and/or physiological variables (e.g., relative humidity,  $\delta_v$ , stomatal 470 conductance). This suggests that environmental and physiological stability is a prerequisite 471 for isotopic steady state to occur. In this context, it should be noted that even when 472 environmental and physiological parameters are stable, the condition of isotopic steady state 473 will not be achieved immediately (e.g. Simonin *et al.* 2013). Rather,  $\delta_E$  will move toward  $\delta_s$ 474 in an exponential manner with a time constant that depends on the leaf water concentration, 475 stomatal conductance, and the water vapour mole fraction inside the leaf (Dongmann et al. 476 1974; Farquhar & Cernusak 2005). Song et al. (2015b) recently conducted a laboratory 477 experiment to monitor this type of exponential trajectory of  $\delta_E$  in cotton leaves exposed to a 478 gas-exchange cuvette environment. They demonstrated that the time constant for the 479 approach of  $\delta_E$  to  $\delta_s$  agreed well with the prediction from the non-steady state isotope theory 480 adapted to cuvette conditions. Under field conditions, time constants for leaf water turnover can often be longer than 481

482 the frequencies at which natural variations in temperature, humidity, and stomatal 483 conductance occur. As a result, it has been argued that the isotopic composition of transpired water,  $\delta_E$ , should rarely be precisely at steady state (Wang & Yakir 1995; Harwood *et al.* 484 485 1998; Simonin *et al.* 2013). A recent field study tracked diurnal variations in  $\delta_E$  for an oak 486 tree during distinct Mediterranean seasons and found that  $\delta_E$  significantly deviated from  $\delta_s$ 487 most of the time (Dubbert et al. 2014). Such an observation, resulting from direct 488 measurements of  $\delta_{\rm E}$ , provides support for the "steady state being rare" argument, thereby 489 suggesting that the steady-state assumption should be used with caution in field conditions 490 when applied to  $\delta_E$ . However, this raises an interesting contrast with  $\delta_L$ , the isotopic 491 composition of leaf water, which often appears to be near to steady state, at least for many  $C_3$ 492 plants, during the day (e.g., Figure 3). This highlights the difference between the isoflux 493  $(E\delta_{\rm E})$  and isostorage  $(W\delta_{\rm L})$  terms, with the latter being relatively buffered against high 494 frequency variations.

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#### 496 Non-steady state effects on leaf water enrichment

497 Non-steady state effects on leaf water isotopic enrichment, the isostorage term, are 498 expected to become important when stomatal conductance is low and/or when leaf water 499 concentrations are high. Most species probably show significant non-steady state behaviour 500 in  $\Delta_{\rm L}$  at night, due to low stomatal conductance (Cernusak, Pate & Farquhar 2002; Cernusak, 501 Farquhar & Pate 2005; Seibt et al. 2006; Barnard et al. 2007; Cuntz et al. 2007). In addition, 502 non-steady state behaviour has been observed during the day in plant species with succulent 503 leaves (Sternberg, Deniro & Johnson 1986; Cernusak et al. 2008), and in some needle-leaved 504 species, when exposed to high vapour pressure deficits or low soil water availability, such 505 that stomatal conductance was relatively low (Pendall, Williams & Leavitt 2005; Seibt et al. 506 2006; Snyder et al. 2010).

507 Variation in leaf water isotopic enrichment under non-steady state conditions ( $\Delta_{Ln}$ ) 508 can be predicted as follows (Farquhar & Cernusak 2005):

509 
$$\Delta_{Ln} = \Delta_L - \frac{\alpha^+ \alpha_k}{g w_i} \cdot \frac{1 - e^{-\wp}}{\wp} \cdot \frac{d(W \Delta_{Ln})}{dt}, \qquad (19)$$

where  $\Delta_L$  is the steady-state prediction of leaf water isotopic enrichment,  $\alpha^+$  is defined as 510  $1+\epsilon^+$ ,  $\alpha_k$  is defined as  $1+\epsilon_k$ , W is the lamina leaf water concentration (mol m<sup>-2</sup>), t is time (s), 511 and g is the total conductance to water vapour of stomata plus boundary layer (mol  $m^{-2} s^{-1}$ ). 512 513 Note that  $\varepsilon^+$  and  $\varepsilon_k$ , if they are expressed in per mil, should be divided by 1000 to calculate  $\alpha^+$ and  $\alpha_k$ . Equation 19 has the term  $\Delta_{Ln}$  on both the left and right sides of the equation, and so 514 515 needs to be solved iteratively. One way to do this is with the Solver function in Microsoft 516 Excel (Farquhar & Cernusak 2005). Alternatively, Kahmen et al. (2008) suggested a simpler 517 method for solving the equation by introducing the assumption that, over sufficiently small time steps,  $\frac{d(W\Delta_{Ln})}{dt} \approx \frac{W\Delta_{Ln} - (W\Delta_{Ln})_{t-1}}{\Delta_t}$ , where the subscript *t*-1 refers to the value at the 518 previous time step and  $\Delta_t$  is the time elapsed since the previous time step. This definition can 519 520 be substituted into Eqn 19, which can then be solved for  $\Delta_{Ln}$ , such that it only occurs on the 521 left side of the equation. The value for  $\Delta_{Ln}$  can then be calculated without need for iteration 522 (Kahmen et al. 2008). Another alternative is to assume a step change in parameters from one 523 time step to the next, so that the leaf water enrichment moves toward the new steady state in 524 an exponential fashion with a time constant,  $\tau$ , approximated by  $W/gw_i$ . This also results in 525 equations that can be calculated without need for iteration (Dongmann et al. 1974; Farquhar 526 & Cernusak 2005; Cuntz et al. 2007).

527

#### 528 Which leaf water model to use?

529 As seen above, models describing leaf water evaporative enrichment have become 530 increasingly complex over time, from the simplest version of the Craig-Gordon equation, to 531 non-steady state models (Dongmann et al. 1974; Farquhar & Cernusak 2005), to the most 532 complex spatially-explicit models describing gradients of enrichment under non-steady state 533 conditions (Cuntz et al. 2007; Ogée et al. 2007). Given the range of options available, it is 534 not always straight forward to decide which leaf water model to use for a particular research 535 question. For some applications at larger temporal and spatial scales, the steady state Craig-536 Gordon model (Eqn 3) will be adequate and including non-steady state effects and Péclet 537 effects will likely add complexity that does not significantly improve model outcomes. 538 Conversely, if water is sampled within a leaf and at high temporal resolution, a spatially-539 explicit and non-steady state model may be required (Ogée et al. 2007). Some studies have tested the suitability of different models in specific applications. For example, Cernusak et 540 541 al. (2005) demonstrated that both the non-steady state and whole-leaf Péclet models were 542 required to predict accurately diel variability in leaf water enrichment in *Eucalyptus globulus*. 543 Ogée et al. (2009) found that the oxygen isotope composition of tree ring cellulose was not 544 sensitive to the value assigned to the Péclet effective length, implying that a simpler two-pool 545 model would have been adequate. At larger spatial but smaller temporal scale, the 546 requirement for a non-steady state model has been confirmed when interpreting variation in 547 ecosystem-scale isofluxes (Xiao et al. 2012; Santos et al. 2014).

Here, we describe a general framework for deciding when to apply different leaf 548 water models. Questions relating to the  $\delta^{18}$ O of oxygen evolution, such as studies of the 549 550 Earth's Dole effect (Bender, Sowers & Labeyrie 1994; Hoffmann et al. 2004), should for the 551 most part be well served by the steady state Craig-Gordon prediction of  $\Delta_{e}$ . This is because oxygen evolution takes place during the day when leaf water is generally near isotopic steady 552 553 state, and because chloroplasts are mostly located near to the evaporative sites. The same argument can be applied for questions relating to effects of photosynthesis on  $\delta^{18}$ O of 554 555 atmospheric CO<sub>2</sub> (Farquhar et al. 1993; Cuntz et al. 2003). However, in this case the impact 556 of exchange of atmospheric  $CO_2$  with leaf water also continues at night. In order to account for the influence of dark respiration on  $\delta^{18}$ O of atmospheric CO<sub>2</sub>, a non-steady state model of 557 558 evaporative site water is needed (Cernusak et al. 2004; Seibt et al. 2006; Cuntz et al. 2007; 559 Santos et al. 2014).

560 For the most part, the influence of leaf water isotopic enrichment on organic material 561 is mediated by photosynthesis. Again, because photosynthesis takes place during the day, research questions relating to  $\delta^{18}$ O and  $\delta^{2}$ H of organic material should be served reasonably 562 563 well by steady state models. Here, there have been mixed results as to whether Péclet effects need to be considered. For isotopic signals closely related to leaf water, such as  $\delta^{18}$ O of 564 565 phloem sugars, a Péclet effect was required (Barbour et al. 2000b; Cernusak et al. 2003). 566 However, in applications that consider processes further downstream from leaf water, such as 567 tree-ring formation, the relatively small Péclet effect becomes further damped, to the point 568 that there may be little advantage in including it (Ogée et al. 2009; Gessler et al. 2014; Song, 569 Clark & Helliker 2014).

570 For plant breeding applications aimed at disentangling effects of stomatal 571 conductance from those of photosynthetic capacity on water-use efficiency, it will likely be 572 advantageous to consider Péclet effects (Farquhar, Condon & Masle 1994; Barbour et al. 573 2000a; Barbour 2007). In addition, Péclet effects will likely be particularly important in 574 studies aimed at linking leaf water stable isotope composition with leaf hydraulic pathways (Barbour & Farquhar 2004; Ferrio et al. 2012; Song et al. 2013). On the other hand, for 575 applications aimed at using the  $\delta^2 H$  of leaf waxes to reconstruct hydrological features of 576 577 ancient ecosystems, the simplest form of the steady state Craig-Gordon equation will likely 578 suffice (McInerney, Helliker & Freeman 2011; Sachse et al. 2012; Kahmen et al. 2013a). 579

#### SAMPLING CONSIDERATIONS AND METHODOLOGICAL ADVANCES 580

581 The isotopic analysis of plant waters presents a number of analytical challenges. 582 These include difficulties of extraction, the necessity to work with small quantities of water, 583 protecting the original composition of the water sample, and avoiding undesirable influences 584 of dissolved compounds. Preventing post-sampling evaporative enrichment of leaf water 585 requires careful consideration of sample handling and storage. For example, even the time taken to separate primary veins from leaf lamina can result in detectable isotopic enrichment 586 587 of the leaf lamina (Cernusak et al. 2003).

588 There are three main ways of analysing the isotopic composition of plant water: 589 equilibration, prior extraction, and simultaneous extraction. With equilibration methods, a 590 gas is equilibrated directly with the plant water while it is still in the sample, and the gas is 591 then analysed for its isotopic composition. With prior extraction, the water is taken out of the plant tissue before isotopic analysis, whereas with simultaneous extraction the water isremoved from the plant as part of the analysis.

594 In equilibration methods, a gas is introduced into a sealed vessel with the sample, 595 such as a detached portion of stem, and the system maintained until the gas has effectively 596 equilibrated with the water that the sample contains (Scrimgeour 1995). Typically, for 597 oxygen isotopic analysis, pure  $CO_2$  or a  $CO_2$ /gas mixture is stored over the sample at a 598 controlled temperature. The  $CO_2$  then exchanges oxygen with the water by the carbonic 599 acid/bicarbonate reaction, with a temperature dependent fractionation. Direct equilibration of 600  $CO_2$  with twig and stem water showed good agreement, to within 0.5‰, with assessments of the  $\delta^{18}$ O of paired samples based on prior extraction (Scrimgeour 1995). This direct 601 equilibration method may also be useful for analysis of the  $\delta^{18}$ O of leaf water. However, a 602 603 limitation may be imposed by the very low rate of diffusion of  $CO_2$  in water, and of water in 604 water, so that the gas may primarily equilibrate with the more exposed portion of the leaf water. Thus, the  $\delta^{18}$ O of the equilibrated CO<sub>2</sub> may be more representative of the evaporative 605 606 site water, as opposed to the bulk leaf water; experiments are needed to test this.

607 Extraction methods aim for complete removal of the water from the sample, because 608 the removal of water by evaporation is typically associated with a fractionation. Thus, in the 609 event of a partial extraction, the water removed will have a different isotopic composition 610 from that which remains. The most widely used prior extraction method is cryogenic vacuum 611 extraction (Ehleringer, Roden & Dawson 2000). Here, the water is freed from the sample 612 using heat and vacuum, and then frozen onto a collecting surface. For plant tissues, 613 cryogenic vacuum extraction is a tested and reliable method, and it typically serves as the 614 benchmark against which new methods are evaluated. However, it is relatively labour and 615 time-intensive. Several authors have proposed modifications aimed at reducing these 616 restrictions (West, Patrickson & Ehleringer 2006a; Vendramini & Sternberg 2007; Koeniger 617 et al. 2011; Ignatev et al. 2013; Orlowski et al. 2013). For laser-based analysis of water 618 isotopes, cryogenic extraction also presents the challenge of transferring organic 619 contaminants that can mix with the water sample and cause optical interference (West et al. 620 2010).

Recent years have seen the advent of laser-based, optical analysers with the capacity
to measure the stable isotope composition of water vapour (e.g. Gupta *et al.* 2009; Sturm &
Knohl 2010; Aemisegger *et al.* 2012; Griffis 2013). Using this type of analyser, new
methods have been developed for simultaneous water extraction and analysis. Here, the leaf

625 is placed in the extraction device and the resulting water vapour is analysed as it is driven off 626 by heating. One example of such a system is an induction module cavity ring down 627 spectroscopy system (IM-CRDS) (Berkelhammer *et al.* 2013). The laser isotope analyser 628 relies on the absorption of an infrared laser pulse by water vapour as it reflects inside a 629 chamber. For a typical liquid water injection, about 1  $\mu$ L of water is vaporised when it is 630 injected into a chamber hotter than boiling point. The vapour is then carried into the analyser 631 in a non-interfering gas. The IM-CRDS system is similar, except that the leaf sample is 632 heated inductively, and the vapour produced is then carried into the laser analyser. A second 633 example of a simultaneous water extraction and analysis system uses a microwave oven to 634 heat the leaf sample (Munksgaard et al. 2014). This was termed ME-IRIS, for microwave 635 extraction isotope ratio infrared spectroscopy. The ME-IRIS system includes a microwave 636 and a condenser to moderate the water vapour concentration of air passing to the laser 637 analyser, so that it remains within the optimal measuring range. Advantages of ME-IRIS are 638 that it can handle larger samples (e.g., whole leaves), and that it uses relatively low cost 639 components, such as a domestic microwave oven.

640 A complication in these simultaneous extraction methods is that some organic 641 compounds, for example alcohols, which can be present in leaf water, interfere significantly 642 with absorption peaks for the target isotopologues in the laser analyser. Two solutions have 643 been developed: a small furnace in-line which breaks down the interfering compounds, and 644 post-processing software that detects and flags analyses that potentially contain spectral 645 interference. The combination of the two tools together appears sufficient to identify and/or 646 reduce the analytical errors associated with organic contaminants to acceptable levels (West 647 et al. 2011; Munksgaard et al. 2014; Martín-Gómez et al. 2015).

648 Both IM-CRDS and ME-IRIS also suffer from memory effects, often requiring two to 649 three sample analyses to overcome the influence of a previous sample if its isotopic 650 composition was substantially different (Berkelhammer et al. 2013; Munksgaard et al. 2014). 651 The impact of the memory effect can be minimized by arranging the analytical sequence in 652 such a way as to avoid large jumps in isotopic composition between adjacent samples. This 653 also highlights a further disadvantage of simultaneous extraction methods; once analysed, the 654 same sample is not available for re-analysis. Thus, wherever possible, samples should be 655 collected in sufficient replication to overcome memory effects and as back-up in the event 656 that a re-analysis is deemed necessary. The main advantage gained by simultaneous 657 extraction is the capacity to analyse samples in the field at the study site or in a temporary

field laboratory, and to thereby have analytical results in near real time so that they caninform the proceeding sampling strategy and experimental design.

660 An interesting variant of an equilibration method and simultaneous analysis of water 661 vapour stable isotopes by a laser analyser has been applied to soil cores (Wassenaar et al. 662 2008). In this system, a soil core was placed inside a sealed, inflatable plastic bag. The 663 sealed bag was then left to equilibrate the water vapour in the headspace with the liquid water 664 in the soil sample. Following the appropriate equilibration time, the plastic bag was punctured with a needle connected to a piece of tubing feeding directly into a laser analyser. 665 The  $\delta^{18}$ O and  $\delta^2$ H of the liquid water in the soil sample could then be inferred from the 666 667 temperature dependent equilibrium fractionation between liquid and vapour,  $\varepsilon^+$ . Such a 668 system may also be suitable for plant materials. Advantages would be the simplicity of the 669 equilibration compared to liquid water extraction, and that isotope ratios of both oxygen and 670 hydrogen could be determined simultaneously. As with the direct equilibration of  $CO_2$ , a 671 question that would need to be addressed for leaves is whether the water vapour in the 672 headspace primarily equilibrates with the evaporative site, or whether it equilibrates with the 673 bulk leaf water.

674

#### 675 CONCLUSIONS

676 Steady state leaf water isotopic enrichment is closely related to relative humidity in 677 natural environments, with the observed enrichment decreasing with increasing relative 678 humidity. Isotopic disequilibrium between source water and atmospheric vapour can also 679 have a relatively strong effect on steady state leaf water isotopic composition. Observations 680 over a large scale environmental gradient in Australia indicated that this effect is likely to be stronger for <sup>2</sup>H than for <sup>18</sup>O. This difference in behaviour between the two isotopes reflects 681 682 the relative magnitudes of the equilibrium and kinetic fractionations in the Craig-Gordon model of evaporative site enrichment. Equilibrium effects dominate for <sup>2</sup>H, whereas kinetic 683 effects dominate for <sup>18</sup>O. 684

In a combined dataset including 118 species, we found that observed bulk leaf water was less enriched than the Craig-Gordon predictions for both <sup>18</sup>O and <sup>2</sup>H, as has been shown previously. Across the full dataset, the proportional difference between Craig-Gordon predicted and observed bulk leaf water enrichment showed no relationship with transpiration rate. Explaining why Péclet effects are detectable in some situations, but not in others, remains a challenge. Linking observed patterns of leaf water isotopic enrichment with specific hydraulic characteristics could provide a tractable way forward, especially withrespect to pathways for water movement from veins to evaporative sites.

693 The development of new technologies for quantifying stable isotope ratios of 694 transpired water and water extracted from plant tissues offers an opportunity to further our 695 understanding of the finer scale controls over leaf water stable isotope enrichment. For 696 example, measuring the isotopic composition of transpired water vapour provides a means of 697 detecting nuances of steady versus non-steady state behaviour, and it also has potential to 698 provide insight into whether slow turnover pools exist within the leaf water, indicative of 699 hydraulic compartmentalisation. Improving our understanding of the environmental and 700 physiological controls over leaf water stable isotopic enrichment will benefit the many 701 applications to which models of this process can be applied, and may additionally lead to novel insights into hydraulic design and functioning in leaves of terrestrial plants. 702

703

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$\Delta_{\rm e}$	Isotopic enrichment of evaporative site water compared to source water
$\Delta_{\rm L}$	Isotopic enrichment of bulk leaf water compared to source water
$\Delta_{Ln}$	Predicted non-steady state isotopic enrichment of bulk leaf water
$\Delta_{ m v}$	Isotopic enrichment of vapour compared to source water (typically negative)
$\delta_{\rm E}$	$\delta^{18}$ O or $\delta^{2}$ H of transpired water vapour
$\delta_{\rm L}$	$\delta^{18}$ O or $\delta^{2}$ H of bulk leaf water
$\delta_{e}$	$\delta^{18}$ O or $\delta^2$ H of water at the evaporative sites within leaves
δs	$\delta^{18}$ O or $\delta^2$ H of source water
$\frac{\delta_v}{\delta^{18}O}$	$\delta^{18}$ O or $\delta^2$ H of atmospheric vapour
$\delta^{18}O$	<sup>18</sup> O/ <sup>16</sup> O relative to the value of a standard (VSMOW for plant waters)
$\delta^2 H$	$^{2}\text{H}/^{1}\text{H}$ relative to the value of a standard (VSMOW for plant waters)
$\epsilon^+$	Equilibrium isotope fractionation between liquid water and vapour
$\epsilon_k$	Kinetic isotope fractionation caused by diffusion of water vapour in air
<i>§</i> )	The Péclet number (representing the ratio between advection and diffusion)
60 E	Transpiration rate
L	Effective path length for water movement through the mesophyll
W	Leaf water concentration
8	Stomatal conductance to water vapour
Wa	Water vapour mole fraction in the atmosphere
Wi	Water vapour mole fraction in the intercellular air spaces inside leaves

Table 1. Symbols and abbreviations used in the text.



**Figure 1.** Spatial distribution of hydrogen and oxygen isotope ratios of precipitation over land. These precipitation isoscapes were derived from a long-term, global network of observations (Welker 2000; IAEA/WMO 2011) and a geostatistical and regression-based model developed with the online workspace IsoMAP (Isoscapes Modeling, Analysis and Prediction v 1.0; <u>http://isomap.org</u>). Model structure, statistical results, and isoscapes may be accessed or downloaded from IsoMAP by referencing job keys: 48170, 48171, 48236, and 48560 (West, 2015).



**Figure 2.** Diel variation in leaf, phloem, pod wall, seed, and xylem water  $\delta^{18}$ O (A) and  $\delta^{2}$ H (B). Samples were collected in Western Australia from *Lupinus angustifolius* grown as part of an agricultural trial. Phloem sap was sampled from pod tips, using a phloem bleeding technique. Redrawn from Cernusak et al. (2002).



**Figure 3.** Relationships between observed leaf water stable isotope enrichment for oxygen  $(\Delta^{18}O_L)$  and hydrogen  $(\Delta^{2}H_L)$  and the relative humidity of the air recorded at the time of sampling (A and C) and the Craig-Gordon predicted enrichments (B and D). Craig-Gordon predicted enrichments were calculated with Eqn 4 of the main text. Samples were collected during daytime from various *Eucalyptus* and *Acacia* species distributed over a sub-continental rainfall gradient in northern Australia (Kahmen *et al.* 2013a). The full dataset is provided in the Supplementary Material. Panel D is redrawn from Kahmen et al. (2013a).



**Figure 4.** Histograms showing the proportional difference between Craig-Gordon predicted leaf water stable isotope enrichment ( $\Delta^{18}O_e$  and  $\Delta^2H_e$ ) and observed bulk leaf water enrichment ( $\Delta^{18}O_L$  and  $\Delta^2H_L$ ) for  $\Delta^{18}O$  (A) and  $\Delta^2H$  (B). Craig-Gordon predicted enrichments were calculated with Eqn 4 of the main text. The dataset combines observations from several publications (Wang, Yakir & Avishai 1998; Cernusak *et al.* 2002; Cernusak *et al.* 2005; Kahmen *et al.* 2008; Kahmen *et al.* 2011; Kahmen *et al.* 2013a; Song *et al.* 2013; Song *et al.* 2014). The full dataset is provided in the Supplementary Material.



**Figure 5.** A schematic representation of the Péclet model of leaf water stable isotopic enrichment (Farquhar & Lloyd 1993). The model describes the average lamina leaf water <sup>18</sup>O and <sup>2</sup>H enrichment relative to that at the evaporative sites as a function of the interplay between diffusion of isotopically enriched water away from the evaporative sites and advection of unenriched vein water toward the evaporative sites. The vein water is transported along a path, the length of which varies as a function of its tortuosity. The average lamina leaf water <sup>18</sup>O and <sup>2</sup>H enrichment decreases if the transpiration rate is high, or when the scaled effective path length is long. Either of these conditions will impede the diffusion of isotopically enriched water away from the evaporative sites. In the figure, red represents the highest <sup>18</sup>O and <sup>2</sup>H enrichment, yellow intermediate, and blue the lowest. The thickness of arrows indicates transpiration rates and the sinuosity of arrows indicates scaled effective path lengths. Redrawn from Cernusak and Kahmen (2013).



**Figure 6.** Schematic of three scenarios for leaf hydraulic design describing the hydraulic linkages between different tissues. The dark grey circle in the middle is a water-filled vein, solid lines depict water flow, with the thicker lines corresponding to higher flow, dashed lines describe diffusion of water vapour, and  $\emptyset$  denotes high resistance between tissue types. In Design 1, the vein is relatively isolated hydraulically from the rest of the leaf; in Design 2, the epidermal tissues are hydraulically linked to the vein by the bundle sheath extensions, but the mesophyll remains relatively isolated; and in Design 3, all tissues are equally well linked hydraulically. Reprinted from Zwieniecki et al. (2007).



**Figure 7.** Spatial variation in leaf water isotopic composition in a tree leaf, *Eucalyptus pauciflora* (A), a banana leaf, *Musa* sp. (B), a cactus stem, *Carnegiea gigantea* (C), and a grass blade, *Miscanthus sinensis* (D). Progressive isotopic enrichment from the base to the apex of the leaf/stem and from the middle toward the edges of the leaf/stem is a common feature. Note that in the banana leaf, the progressive enrichment shows in the perpendicular direction to the midrib, rather than along its length. The left side of the banana leaf shows the vein patterning. Scale bars are approximate. Figures are modified from Santrucek *et al.* (2007), Stuart-Williams (unpublished), English *et al.* (2007), and Helliker and Ehleringer (2000), respectively.



### Farquhar-Gan Péclet Model

**Figure 8.** A schematic representation of the Farquhar and Gan (2003) Péclet model that predicts progressive enrichment of leaf water stable isotopes along a leaf. Darker blue indicates a higher level of stable isotope enrichment. In the model, isotopically lighter water is preferentially transpired leaving heavier water to diffuse back into the xylem and be carried further along the leaf. For this pattern to be pronounced, the ratio of advection to diffusion (Péclet number) has to be large in the longitudinal direction, and small in the radial direction. In the figure,  $\wp_1$  is the longitudinal Péclet number,  $\wp_{rv}$  is the radial Péclet number associated with veinlets, and  $\wp$  is the radial Péclet number associated with the mesophyll in the leaf lamina.